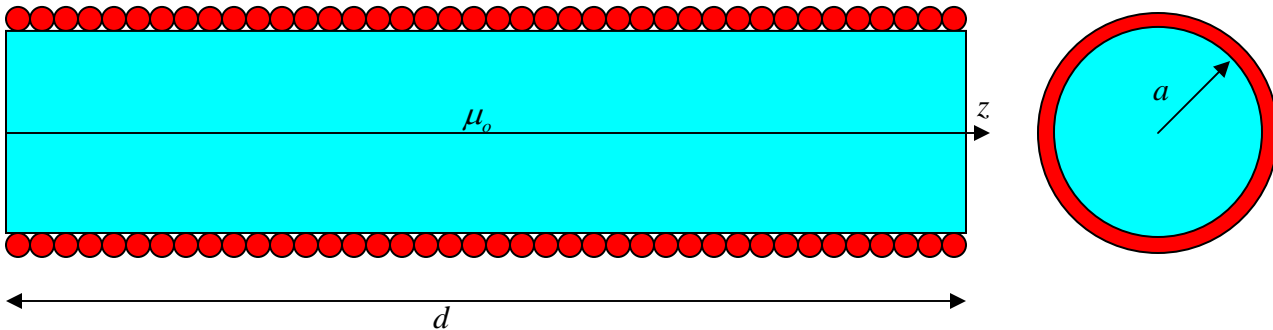




1. Maxwell's Equations – Quasi-statics



An air core, N turn, cylindrical solenoid of length d and radius a , carries a current $I = I_0 \cos \omega t$.

- Using Ampere's Law, determine the magnetic flux density \vec{B} and magnetic field intensity \vec{H} for this solenoid, assuming that the solenoid is very long and there is no fringing. Note that, while you may know the answer to this question by inspection, show all steps. That is, set up and evaluate the line integral, find the current enclosed by the line, etc.
- Evaluate the time derivative of the magnetic flux density $\frac{\partial \vec{B}}{\partial t}$.
- From the point form of Faraday's Law, we know that $\nabla \times \vec{E} = \frac{\partial \vec{B}}{\partial t}$ or

$$\hat{r} \left(\frac{1}{r} \frac{\partial \mathcal{E}_z}{\partial \phi} - \frac{\partial \mathcal{E}_\phi}{\partial z} \right) + \hat{\phi} \left(\frac{\partial \mathcal{E}_r}{\partial z} - \frac{\partial \mathcal{E}_z}{\partial r} \right) + \hat{z} \frac{1}{r} \left(\frac{\partial (r \mathcal{E}_\phi)}{\partial r} - \frac{\partial \mathcal{E}_r}{\partial \phi} \right) = \hat{z} \frac{\partial \mathcal{B}_z}{\partial t} \text{ where we have}$$

used the fact that the magnetic flux density $\vec{B} = \hat{z} B_z(r)$ to drop the x and y directed terms from the right hand side. We can use this information to also drop

the same terms from the left hand side. $\hat{z} \frac{1}{r} \left(\frac{\partial (r \mathcal{E}_\phi)}{\partial r} - \frac{\partial \mathcal{E}_r}{\partial \phi} \right) = \hat{z} \frac{\partial \mathcal{B}_z}{\partial t}$. So far, we

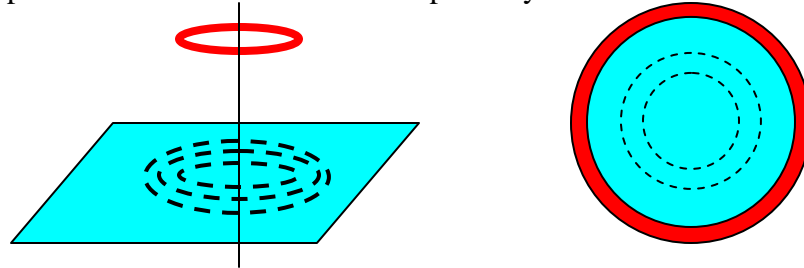
have shown that the electric field \vec{E} can have, at most a ϕ or r directed component. In fact, $\vec{E} = \hat{\phi} \mathcal{E}_\phi(r)$ only because the electric field in this case has to be in the same direction as the currents that produce the solenoid \vec{B} field. Why this must be is a bit subtle, but maybe best seen by looking at what would happen if the core material was iron instead of air. If that was the case, then we would observe eddy currents being created in the iron that mimic the current in the solenoid winding in an attempt to reduce the applied field as much as possible.



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Recall that currents induced in the copper pipe or copper plate by the strong permanent magnet (demo done in class) we in a direction to produce a field in the opposite direction. To qualitative show this for the solenoid, the end view of the solenoid has been reproduced below with some example eddy currents shown as dashed lines.

Also shown is a generic circular current inducing eddies in a flat plate.



Now, since there is only one component for $\vec{E} = \hat{\phi}E_{\phi}(r)$, Faraday's Law

simplifies to $\frac{1}{r} \left(\frac{\partial(rE_{\phi})}{\partial t} \right) = \frac{\partial B_z}{\partial t}$. Solve this equation for $\vec{E} = \hat{\phi}E_{\phi}(r)$. Note that

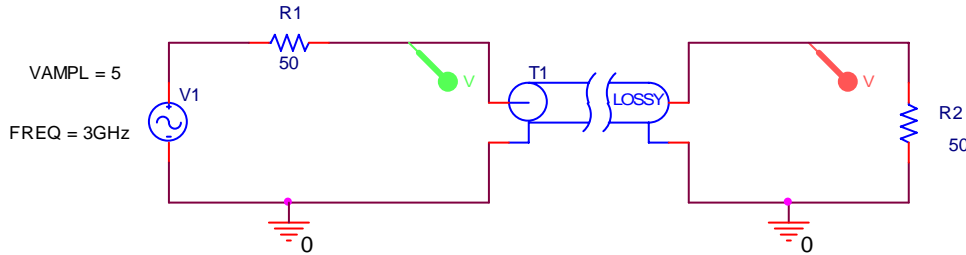
your solution must be a function of radius r and cannot be a constant like \vec{B} since \vec{E} must have a curl.

- d. The solutions to parts a and c are the quasi-static field. However, to determine the accuracy of this solution, we need to use the \vec{E} just found to find a correction to \vec{B} . This correction to the original field must also be a function of r . To start this process, determine the electric flux density \vec{D} from \vec{E} and then find its time derivative $\frac{\partial \vec{D}}{\partial t}$. This is the displacement current density.
- e. Use Ampere's Law, following the same steps as in part a, to find the new contribution to the magnetic field intensity \vec{H} . Note that, since the displacement current density is distributed throughout the region $0 \leq r \leq a$, evaluating the right hand side of Ampere's Law is a bit more complicated than it was for part a. Once you have the new \vec{H} , also find the new \vec{B} .
- f. At this point, we can represent our field solutions as $\vec{B} = \vec{B}_0 + \vec{B}_2$ and $\vec{E} = \vec{E}_1$ where the subscripts indicate the order of the corrections. For quasi-statics, we should have that $\frac{B_2}{B_0} \ll 1$. Evaluate this ratio and determine the range of frequencies for which the quasi-static solution (i.e. $\vec{B} = \vec{B}_0$ and $\vec{E} = \vec{E}_1$) is reasonably accurate.

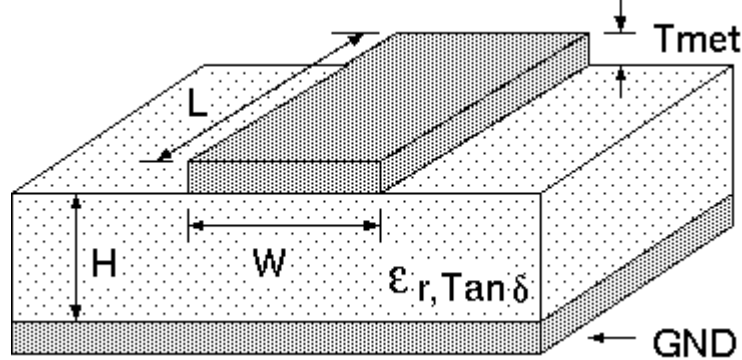
Hint: There is an extensive discussion of quasi-statics for two other geometries in the class notes, Unit 9, pages 17-24.



2. Plane Waves and Transmission Lines



A transmission line is driven by a 5V, 3GHz voltage source. Both the source and load resistances are 50 Ohms. The length of the line is 10cm. Now that we have addressed finding capacitance and inductance, we should begin such a problem by first describing the configuration of the transmission line. Assume that we have a microstrip line for which there are many, many tools available to determine the line characteristics. The parameters of the line are as follows: $\epsilon_r = 4.2$, dielectric thickness = 0.79375mm, the trace width is 1.5mm, the trace thickness is 0.035mm. A view of the line is shown below. Note that this shows a loss tangent, but we will assume lossless conditions.

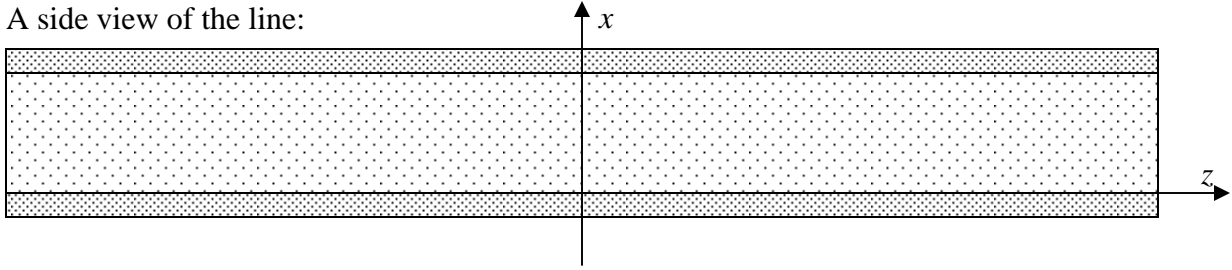


- Using this information and one of the impedance calculators found on the Resources webpage, determine the characteristic impedance Z_o , the delay time t_d , the capacitance and inductance per unit length for this line. You should find that the line is matched to the source and load.
- Use your answers to part a to set up the PSpice simulation shown above. Do time-domain analysis and show roughly 3-5 periods of the source and load voltages.
- Write the phasor expressions for the voltage and current on the line and verify that your expression is consistent with your PSpice solution.

Now, will use the information have on voltage and current to find the electric and magnetic fields that exist in the region between the plates. For this analysis, we will assume that the plates are infinite so that there is no fringing and the field structure is that of an ideal parallel plate. This approach will give us a good answer for the middle of the field region, but probably not near the edges of the plates. We will first begin by finding the electric field. We will not find the magnetic field directly. Rather, we will find it from the electric field.



A side view of the line:



- d. Assuming that the phasor voltage on the line creates an electric field in the insulator region between the plates, find the electric field \vec{E} . Note the coordinate system shown above and that this field is a uniform plane wave.
- e. From \vec{E} find \vec{D} and then determine the surface charge density ρ_s on each conductor using the boundary condition on the normal component of \vec{D} .
- f. Determine the intrinsic impedance of the dielectric medium η .
- g. From your answers for \vec{E} and η , determine the magnetic field intensity \vec{H} .
- h. Show that your answer for \vec{H} is consistent with your expression for the current on the line (from part c) by evaluating the boundary condition for the tangential component of \vec{H} .
- i. Using your expressions for \vec{E} and \vec{H} , determine the average power density (Poynting vector) and then find the total power flowing in the line (see below).
- j. Show that the total power flowing in the line is equal to the power delivered to the load resistor from your transmission line analysis.

Note: Because the configuration addressed here is not an ideal parallel plate structure, we would have to solve for the fields using a numerical method to answer this question directly. However, we can address an equivalent ideal structure to see that the voltages and fields are properly connected. To answer parts h, i and j, use the expressions in table 2.1 of Ulaby for inductance and capacitance per unit length for a parallel plate structure. Adjust the values for ϵ and w (keeping the value of d the same) to obtain the identical values for Z_o , l , and c for the transmission line. Then answer the questions.